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# Anomalous magnetic scattering in tetragonal $\text{RB}_2\text{C}_2$ ( $\text{R} = \text{rare earth}$ ) observed by means of neutron diffraction

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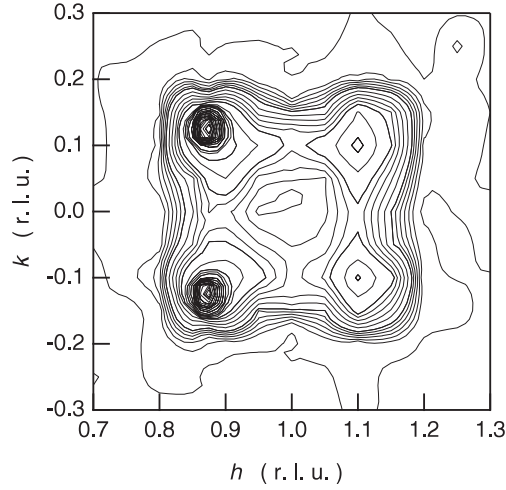
## Abstract

Detailed neutron diffraction experiments have been performed on tetragonal  $\text{R}^{11}\text{B}_2\text{C}_2$  ( $\text{R} = \text{rare-earth}$ ) compounds which show characteristic antiferroquadrupolar orderings.  $\text{TbB}_2\text{C}_2$ ,  $\text{ErB}_2\text{C}_2$  and  $\text{HoB}_2\text{C}_2$ , which have magnetic long-periodic states, show anomalous magnetic diffuse scattering with flat intensity. The anomalous diffuse scattering exists in a small region in the reciprocal space, surrounded by satellite peaks, due to the long-periodic magnetic structures. The profile and temperature dependence of the anomalous diffuse scattering indicate that the origin cannot be understood on the basis of magnetic short-range correlations.

## 1. Introduction

$\text{RB}_2\text{C}_2$  ( $\text{R}$ : rare-earth) compounds, which have the tetragonal  $\text{LaB}_2\text{C}_2$ -type crystal structure [1–3], show diverse and anomalous magnetic and electric quadrupolar behaviours caused by competition between antiferromagnetic (AFM) and antiferroquadrupolar (AFQ) interactions [4, 5]. Of the  $\text{RB}_2\text{C}_2$  system,  $\text{TbB}_2\text{C}_2$  and  $\text{HoB}_2\text{C}_2$  show particularly characteristic AFQ orderings.  $\text{HoB}_2\text{C}_2$  exhibits an AFQ ordering at  $T_Q = 4.5$  K under zero magnetic field, even though a magnetic ordered state exists between  $T_Q$  and  $T_N = 5.9$  K ( $>T_Q$ ) [5–7]. The AFQ ordering in  $\text{HoB}_2\text{C}_2$  was also confirmed by resonant x-ray scattering experiments by Matsumura *et al* [8]. On the other hand,  $\text{TbB}_2\text{C}_2$  shows an AFM ordering at  $T_N = 21.7$  K, but no AFQ ordering under zero magnetic field has been found so far. However, Kaneko *et al* [9, 10] reported that there exist AFQ ordered states under magnetic fields in  $\text{TbB}_2\text{C}_2$ . Since the internal and external magnetic fields lift the degeneracy of the ground states, the AFQ orderings in  $\text{TbB}_2\text{C}_2$  and  $\text{HoB}_2\text{C}_2$  are unique phenomena.

We think that the magnetic long-periodic states which were found in  $\text{HoB}_2\text{C}_2$  and  $\text{TbB}_2\text{C}_2$  are important for investigation of the anomalies in  $\text{RB}_2\text{C}_2$ . The magnetic ordered state in  $\text{TbB}_2\text{C}_2$  below  $T_N$  has a long-periodic magnetic component with the magnetic moments in



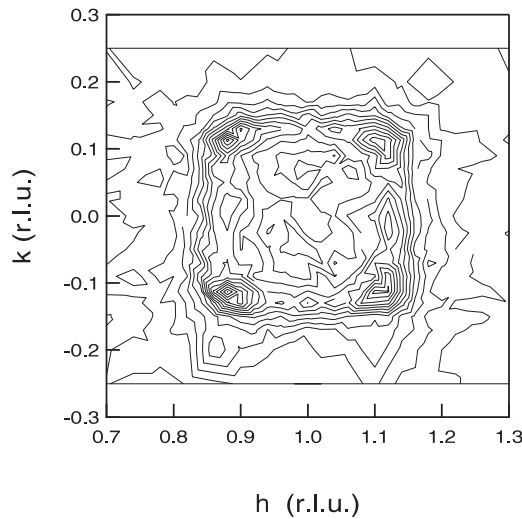
**Figure 1.** A contour map of the anomalous magnetic scattering at  $T = 7$  K around the  $(1, 0, 0)$  position in  $\text{Tb}^{11}\text{B}_2\text{C}_2$  observed for KSD. Details have been reported in [11].

the  $c$ -plane as an appendant component of a very dominant  $\mathbf{k} = (1, 0, 1/2)$  type of AFM structure with a small  $\mathbf{k}' = (0, 0, 1/2)$  component; the propagation vector of the long-periodic component is  $\mathbf{k}_L = (1 + \delta, \delta, 0)$ , where  $\delta = 0.13$  [9]. On the other hand, the long-periodic magnetic phase in  $\text{HoB}_2\text{C}_2$  exists between  $T_N$  and  $T_Q$  [7, 12]. The propagation vector in the intermediate phase is  $\mathbf{k}_L = (1 + \delta, \delta, \delta')$ , where  $\delta = 0.112$  and  $\delta' = 0.04$  [7]. Note that  $\text{ErB}_2\text{C}_2$ , which shows no AFQ ordering, has a long-periodic magnetic state as well; the propagation vector in  $\text{ErB}_2\text{C}_2$  is  $\mathbf{k}_L = (1 + \delta, \delta, 0)$ , where  $\delta = 0.112$  [13]. The close similarity of the periodicities probably indicates that the long periodicity is based on some characteristics common to all members of the  $\text{RB}_2\text{C}_2$  system. Thus, it is important to compare the long-periodic states with each other to understand characteristics of  $\text{RB}_2\text{C}_2$ .

The long-periodic magnetic structure phases in  $\text{HoB}_2\text{C}_2$  and  $\text{TbB}_2\text{C}_2$  are characterized by anomalous broad magnetic scattering with a trapezium-type profile as well as satellite peaks with long-periodic structures. Figure 1 shows the anomalous magnetic diffuse scattering observed in  $\text{TbB}_2\text{C}_2$  at  $T = 7$  K around the  $(1, 0, 0)$  position [11]. The profile of the magnetic diffuse scattering cannot be understood in terms of magnetic short-range correlation or critical phenomena around  $T_N$ . Details of the anomalous scattering in  $\text{TbB}_2\text{C}_2$  have been reported in [11]. The same trapezium-type diffuse scattering was also observed in the long-periodic state in  $\text{HoB}_2\text{C}_2$  [12]. The anomalous magnetic scattering is thought to be an important feature in the  $\text{RB}_2\text{C}_2$  system because of its similarity to those in  $\text{HoB}_2\text{C}_2$  and  $\text{TbB}_2\text{C}_2$ ; however, the origin of the anomalous magnetic scattering, especially the roles of AFQ interactions, is unknown as yet.

In  $\text{ErB}_2\text{C}_2$  which also has a long-periodic magnetic structure phase between  $T_N = 15.9$  K and  $T_l = 13.0$  K, a similar broad magnetic scattering was observed in powder diffraction experiments [13]. Since  $\text{ErB}_2\text{C}_2$  shows no AFQ ordering, comparing the broad magnetic scattering in  $\text{ErB}_2\text{C}_2$  with those in  $\text{HoB}_2\text{C}_2$  and  $\text{TbB}_2\text{C}_2$  is important for understanding the possible role of AFQ interactions in the anomalous scattering. The authors, therefore, performed neutron diffraction experiments on a single-crystalline sample of  $\text{Er}^{11}\text{B}_2\text{C}_2$  to clarify the diffuse scattering observed in the powder diffraction experiments.

For sample preparation, mixtures of 99.9% pure Er, 99.5%-enriched  $^{11}\text{B}$  and 99.999% pure C were melted by the conventional argon arc technique; natural boron was replaced to



**Figure 2.** A contour map of the magnetic scattering at  $T = 15.0$  K ( $\sim T_N$ ) around the  $(1, 0, 0)$  position in the  $(h, k, 0)$  plane in  $\text{Er}^{11}\text{B}_2\text{C}_2$  observed for KSD. There exists a square region surrounded by satellite positions,  $(1 + \delta, \pm\delta, 0)$  and  $(1 - \delta, \pm\delta, 0)$ , with a nearly flat region around  $(1, 0, 0)$ . At this temperature, no satellite peak due to the long-periodic structure was detectable.

avoid its strong absorption effect. The single crystals were grown by the Czochralski method using a tri-arc furnace.

We performed neutron diffraction experiments on the Kinken neutron diffractometer KSD installed at the reactor JRR-3M in Japan Atomic Energy Research Institute, Tokai. A neutron beam with  $\lambda = 1.52$  Å was obtained from the 3 1 1 reflection of the Ge monochromator. The collimation condition was  $12'$ -open-sample- $30'$ . The single-crystalline sample was mounted at the cold head of a closed-cycle He-gas refrigerator with the  $c$ -plane horizontal.

## 2. Results

In the experiments, anomalous broad magnetic scattering was also observed even in  $\text{Er}^{11}\text{B}_2\text{C}_2$ ; the characteristics of the anomalous scattering are basically the same as those observed in the AFQ ordering compounds  $\text{HoB}_2\text{C}_2$  and  $\text{TbB}_2\text{C}_2$ . Figure 2 shows a contour map of the intensity of the magnetic scattering in  $\text{ErB}_2\text{C}_2$  around the  $(1, 0, 0)$  position at  $T = 15.0$  K; no magnetic satellite peak was observed at that temperature. As seen in figure 2, the profile of the magnetic diffuse scattering in  $\text{ErB}_2\text{C}_2$  is nearly the same as that in  $\text{TbB}_2\text{C}_2$  shown in figure 1 except for the existence of satellite peaks for  $\text{TbB}_2\text{C}_2$  at  $T = 7$  K; the nearly flat region was also observed in  $\text{ErB}_2\text{C}_2$ . The flat region indicates that the anomalous magnetic scattering shown in figure 2 cannot be understood in terms of magnetic short-range correlations either. With increasing temperature, anomalous diffuse scattering develops above  $13$  K  $\sim T_i$ , where satellite peaks also develop, and has its intensity maximum at about  $15$  K  $\sim T_N$ , where the satellite peaks disappear. Note that the anomalous diffuse scattering remains up to about  $20$  K.

The important features of the magnetic long-periodic states in  $\text{RB}_2\text{C}_2$  (R = Tb, Ho and Er) which were found by our neutron diffraction experiments are as follows:

- (1) there exists anomalous magnetic scattering in the square region around the  $(1, 0, 0)$  reciprocal-lattice position with nearly flat intensity surrounded by the magnetic satellite peaks;

- (2) diffuse-type magnetic scattering components were also observed around and between the satellite positions;
- (3) these magnetic scattering components show complicated temperature dependence and persist up to  $\sim 2T_N$  or  $2T_Q$ ;
- (4) with increasing temperature, the intensity of the anomalous magnetic scattering in  $\text{ErB}_2\text{C}_2$  shows a maximum at  $T_N$ , where the satellite peaks due to the long-periodic structure disappear, while those for  $\text{HoB}_2\text{C}_2$  and  $\text{TbB}_2\text{C}_2$  show no obvious anomaly at the temperatures where satellite peaks due to long periodicity disappear.

Details of the characteristics of the anomalous diffuse scattering have been reported in [11, 12].

The fact that  $\text{ErB}_2\text{C}_2$  shows the same type of anomalous magnetic scattering as  $\text{HoB}_2\text{C}_2$  and  $\text{TbB}_2\text{C}_2$  indicates that AFQ interactions are not necessarily required to understand the origin of the anomalous magnetic scattering. Moreover, the anomalous magnetic scattering is a characteristic common to all of the  $\text{RB}_2\text{C}_2$  compounds with long-periodic magnetic structures. From this result and the closely similar periodicity, the authors conclude that characteristics found for the long-periodic magnetic states of  $\text{HoB}_2\text{C}_2$ ,  $\text{TbB}_2\text{C}_2$  and  $\text{ErB}_2\text{C}_2$  are based on a property common to all members of the  $\text{RB}_2\text{C}_2$  system, such as a structure of the Fermi surface and/or a common characteristic of the magnetic interactions in  $\text{RB}_2\text{C}_2$ .

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